

# MAXWELL PLASMA AND QUANTUM FLUCTUATIONS OF ELECTROMAGNETIC FIELD

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Generally they construct the theory of high temperature classical electron –ion plasma being in thermodynamical state using Vlasov’s equation [1,2]. In such a way one gets the plasma oscillations with frequency  $\Omega_e = \sqrt{e^2 n_e / m_e}$ . This frequency does not depend on wave length. The quantum electrodynamics by taken into account the quantum properties of electromagnetic field brings in theory the dimensionless parameter [3]  $Z = e^2 T / 2\pi^2 \hbar^2 v_e \Omega_e$ . This parameter connecting with Planck constant  $\hbar$  dramatically changes the results of the theory. Here  $v_e = \sqrt{T/m}$ . Theoretically the Maxwell plasma may be as in state  $Z \ll 1$  (without the quantum fluctuations of electromagnetic field) and in state  $Z \gg 1$  (in present of the quantum fluctuations of electromagnetic fields). The Vlasov’s theory deals with the case  $Z = 0$ . The linear with respect to  $Z$  correction terms to the Vlasov’s theory are inversional proportional to  $\hbar^2$ , This fact excludes as the limit  $\hbar \rightarrow 0$  and the existence of Vlasov’s theory itself. The aggressive feed –back action of thermal exciting electromagnetic Langmuir waves on plasma electrons and ions describing by  $Z$  parameter dramatically changes the dispersion equations. In real conditions in Maxwell plasma the parameter  $Z \gg 1$  is much larger then unity. That is why the properties of plasma dramatically differ of Vlasov’s equation descriptions. Instead of Langmuir’s frequency  $\Omega_e$  one has now an another character quantum frequency  $\omega = \Omega_e \sqrt{2Z} = (1/\hbar\pi) \sqrt{Te^2/r_D}$ , The two branches of electron oscillations appear. One of them describes the electron sound with the velocity  $V_e = v_e \sqrt{\hbar\omega th(\hbar\omega/T)/2T}$ . The ions of plasma possess as the sound branch of spectrum and another branch which rests finite by small wave numbers [4]. The properties of transversal and longitudinal dielectric permittivity’s dramatically differ each to another especially at small wave numbers.

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